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A VERY FAST RAMPING MUON SYNCHROTRON FOR A NEUTRINO FACTORY

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Abstract

A 4600 Hz fast ramping synchrotron is studied as an economical way of accelerating muons from 4 to 20 GeV/c for a neutrino factory. Eddy current losses are minimized by the low machine duty cycle plus thin grain oriented silicon steel laminations and thin copper wires. Combined function magnets with high gradients alternating within single magnets form the lattice. Muon survival is 83%.

Historically synchrotrons have provided economical particle acceleration. Here we explore a very fast ramping muon synchrotron [1] for a neutrino factory [2]. The accelerated muons are stored in a racetrack to produce neutrino beams ($\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$ or $\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu$). Neutrino oscillations [3] have been observed at experiments [4] such as Homestake, Super-Kamiokande, SNO, and KamLAND. Further exploration using a neutrino factory could reveal CP violation in the lepton sector.

This synchrotron must accelerate muons from 4 to 20 GeV/c with moderate decay loss ($\tau_{\mu^\pm} = 2.2 \mu\text{S}$). Because synchrotron radiation goes as m^4 , muons radiate two billion times ($(105.7/.511)^4$) less power than electrons for a given ring diameter and lepton energy. Grain oriented silicon steel (3% Si) is used to provide a high magnetic field with a high μ to minimize magnetic energy ($B^2/2\mu$) stored in the yoke. Magnetic energy stored in the gap is minimized by reducing its size. Cool muons [5] with low beam emittance allow this. The voltage needed to drive a magnet is $-L di/dt$. Voltage is minimized by shrinking the volume of stored energy to reduce the inductance, L .

Acceleration to 4 GeV/c might feature fixed field dog-bone arcs [6] to minimize muon decay loss. Fast ramping synchrotrons [6, 7] might also accelerate cold muons to higher energies for a $\mu^+ \mu^-$ collider [8].

We form arcs with sequences of combined function cells within continuous long magnets, whose poles are alternately shaped to give focusing gradients of each sign. A cell has been simulated using SYNCH [9]. Gradients alternate from positive 20 T/m gradient (2.24 m long), to zero gradient (.4 m long) to negative 20 T/m gradient (2.24 m) to zero gradient (0.4 m), etc. Details are given in Table 1.

It is proposed to use 5 such arc cells to form an arc segment. These segments are alternated with straight sections containing RF. The phase advance through one arc segment is $5 \times 72^\circ = 360^\circ$. This being so, dispersion suppression between straights and arcs can be omitted. With no dispersion

Table 1: Combined function magnet cell parameters. Five cells make up an arc and 18 arcs form the ring.

Cell length	m	5.28
Combined Dipole length	m	2.24
Combined Dipole B_{central}	Tesla	0.9
Combined Dipole Gradient	T/m	20.2
Pure Dipole Length	m	0.4
Pure Dipole B	Tesla	1.8
Momentum	GeV/c	20
Phase advance/cell	deg	72
beta max	m	8.1
Dispersion max	m	0.392
Normalized Trans. Acceptance	π mm rad	4

in the straight sections, the dispersion performs one full oscillation in each arc segment, returning to zero for the next straight as shown in Fig. 1. There will be 18 such arc segments and 18 straight sections, forming 18 superperiods in the ring. Straight sections (22 m) without dispersion are used for superconducting RF, and, in two longer straights (44 m), the injection and extraction. To assure sufficiently low magnetic fields at the cavities, relatively long field free regions are desirable. A straight consisting of two half cells would allow a central gap of 10 m between quadrupoles, and two smaller gaps at the ends. Details are given in Table 2. Matching between the arcs and straights is not yet designed. The total circumference of the ring including combined functions magnets and straight sections adds up to 917 m ($18 \times 26.5 + 16 \times 22 + 2 \times 44$).

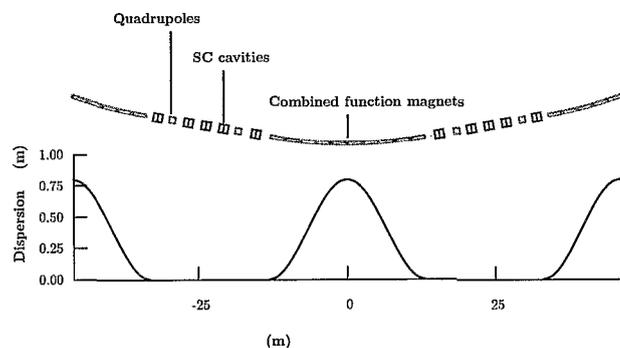


Figure 1: Combined function magnets bend the muons in the arcs. Superconducting RF cavities accelerate muons in the straight sections. Two quadrupoles per straight section provide focusing. The straight sections are dispersion free.

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Table 2: Straight section lattice parameters.

ϕ	$L_{\text{cell}}/2$	L_{quad}	dB/dx	a
77°	11 m	1 m	7.54 T/m	5.8 cm
β_{max}	σ_{max}	B_{pole}	$U_{\text{mag/quad}}$	
36.6 m	.0195 m	0.44 T	≈ 3000 J	

The RF (see Table 3) must be distributed around the ring to avoid large differences between the beam momentum (which increases in steps at each RF section) and the arc magnetic field (which increases continuously). The amount of RF used is a tradeoff between cost and muon survival. Survival is somewhat insensitive to the fraction of stored energy the beam removes from the RF cavities, because the voltage drop is offset by time dilation.

The muons accelerate from 4 to 20 GeV. If they are extracted at 95% of full field they will be injected at 19% of full field. For acceleration with a plain sine wave, injection occurs at 11° and extraction occurs at 72° . So the phase must change by 61° in 37 μsec . Thus the sine wave goes through 360° in 218 μsec , giving 4600 Hz.

Estimate the energy stored in each 26.5 m long combined function magnet. The gap is about .14 m wide and has an average height of .06 m. Assume an average field of 1.1 Tesla. The permeability constant, μ_0 , is $4\pi \times 10^{-7}$. $W = B^2/2\mu_0[\text{Volume}] = 110\,000$ Joules. Next given one turn, an LC circuit capacitor, and a 4600 Hz frequency; estimate current, voltage, inductance, and capacitance.

$$B = \frac{\mu_0 NI}{h} \rightarrow I = \frac{Bh}{\mu_0 N} = 52 \text{ kA} \quad (1)$$

$$W = .5 LI^2 \rightarrow L = 2W/I^2 = 80 \mu\text{H} \quad (2)$$

$$f = \frac{1}{2\pi} \sqrt{\frac{1}{LC}} \rightarrow C = \frac{1}{L(2\pi f)^2} = 15 \mu\text{F} \quad (3)$$

$$W = .5 CV^2 \rightarrow V = \sqrt{2W/C} = 120 \text{ kV} \quad (4)$$

Table 3: Superconducting RF parameters.

Frequency	201	MHz
Gap	.75	m
Gradient	15	MV/m
Stored Energy	900	Joules
Muons per train	5×10^{12}	
Orbits (4 to 20 GeV/c)	12	
No. of RF Cavities	160	
RF Total	1800	MV
ΔU_{beam}	110	Joules
Energy Loading	.082	
Voltage Drop	.041	
Muon Acceleration Time	37	μsec
Muon Survival	.83	

Separate coils might be put around each return yoke to halve the voltage. The stack of SCRs driving each coil might be center tapped to halve the voltage again. Four equally spaced 6 cm coil slots could be created in each side yoke using 6 cm of wider laminations to cut the voltage by five, while leaving the pole faces continuous. 6 kV is easier to insulate than 120 kV. It may be useful to shield [1] or chamfer [10] magnet ends to avoid large eddy currents where the field lines typically do not follow laminations. A DC offset power supply could be useful [1]. Neutrino horn power supplies look promising.

Grain oriented silicon steel is chosen for the return yoke due to its high permeability at high field as noted in Table 4. The skin depth [11], δ , of a lamination is 160 μm from eqn. 5. ρ = resistivity.

$$\delta = \sqrt{\rho / \pi f \mu} = \sqrt{47 \times 10^{-8} / \pi 4600 1000 \mu_0} \quad (5)$$

Take $\mu = 1000\mu_0$ as a limit on magnetic saturation and hence energy storage in the yoke. Next estimate the fraction of the inductance of the yoke that remains after eddy currents shield the laminations [12]. The lamination thickness, t , equals 100 μm [13]. $L/L_0 = (\delta/t) (\sinh(t/\delta) + \sin(t/\delta)) / (\cosh(t/\delta) + \cos(t/\delta)) = 0.995$. So it appears that magnetic fields can penetrate 100 μm thick laminations at 4600 Hz. Thicker 175 μm thick laminations [14] would be half as costly and can achieve a somewhat higher packing fraction. $L/L_0(t = 175 \mu\text{m}) = 0.956$.

Calculate the resistive energy loss in the copper coils. There are four 5 cm square copper conductors each 5300 cm long. $R = 5300 (1.8 \mu\Omega\text{-cm}) / (4) (5^2) = 95 \mu\Omega$. So, $P = I^2 R \int_0^{2\pi} \cos^2(\theta) d\theta = 130\,000$ w/magnet. Eighteen magnets give a total loss of 2340 kw. But the neutrino factory runs at 30 Hz. Thirty half cycles of 109 μsec per second gives a duty factor of 300 and a total $I^2 R$ loss of 8000 watts. Muons are orbited in opposite directions on alternate cycles. If this proves too cumbersome, the duty cycle factor could be lowered to 150. See if .25 mm (30 gauge) wire is usable. From eqn. 6, the skin depth, δ , of copper at 4600 Hz is 0.97 mm.

$$\delta = \sqrt{\rho / \pi f \mu_0} = \sqrt{1.8 \times 10^{-8} / \pi 4600 \mu_0} \quad (6)$$

Now calculate the dissipation due to eddy currents in this $w = .25$ mm wide conductor, which will consist of

Table 4: Permeability ($B/\mu_0 H$). Grain oriented silicon steel has a far higher permeability parallel (\parallel) to than perpendicular (\perp) to its rolling direction [14, 15]. T = Tesla.

Material	1.0 T	1.5 T	1.8 T
1008 Steel	3000	2000	200
Grain Oriented (\parallel)	40000	30000	3000
Grain Oriented (\perp)	4000	1000	

transposed strands to reduce this loss [10, 16]. To get an idea, take the maximum B-field during a cycle to be that generated by a 0.025m radius conductor carrying 26000 amps. The eddy current loss in a rectangular conductor made of transposed square wires .25 mm wide (Litz wire [17]) with a perpendicular magnetic field is as follows. $B = \mu_0 I / 2\pi r = 0.2$ Tesla.

$$P = [\text{Volume}] \frac{(2\pi f B w)^2}{24\rho} \quad (7)$$

$$P = [4 .05^2 53] \frac{(2\pi 4600 .2 .00025)^2}{(24) 1.8 \times 10^{-8}} = 2800 \text{ kw} \quad (8)$$

Multiply by 18 magnets and divide by a duty factor of 300 to get an eddy current loss in the copper of 170 kw. Stainless steel water cooling tubes will dissipate a similar amount of power [6]. Alloy titanium cooling tubes would dissipate half as much.

Do the eddy current losses [16] in the 100 μm thick iron laminations. Use eqn. 7 with a quarter meter square area, a 26.5 m length, and an average field of 1.1 Tesla.

$$P = [(26.5) (.5^2)] \frac{(2\pi 2600 1.1 .0001)^2}{(24) 47 \times 10^{-8}} = 5900 \text{ kw} \quad (9)$$

Multiply by 18 magnets and divide by a duty factor of 300 to get an eddy current loss in the iron laminations of 350 kw or 700 watts/m of magnet. So the iron will need some cooling. The ring only ramps 30 time per second, so the $\int \mathbf{H} \cdot d\mathbf{B}$ hysteresis losses will be low, even more so because of the low coercive force ($H_c = 0.1$ Oersteds) of grain oriented silicon steel. This value of H_c is eight times less than 1008 low carbon steel.

The low duty cycle of the neutrino factory leads to eddy current losses of less than a megawatt in a 4600 Hz, 917 m circumference ring. Muon survival is 83%. The high permeability of grain oriented silicon steel permits high fields with little energy stored in the yoke. Gradients are switched within dipoles to minimize eddy current losses in ends. Time dilation allows extra orbits with little muon decay at the end of a cooling cycle. This allows one to use more of the stored RF energy. Much of the magnetic field in our lattice is used for focusing rather than bending the muon beam. More muon cooling would lead to less focusing, more bending, and an even smaller ring.

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